

Physics 403: Relativity

Homework Assignment 2

Due 12 February 2007

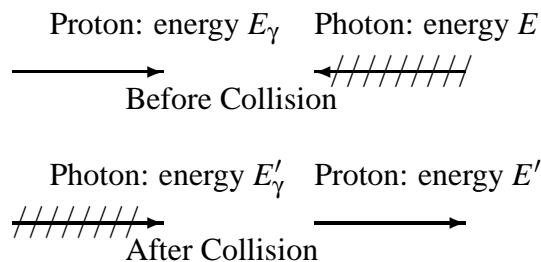
1. Inverse Compton scattering occurs whenever a photon scatters off a particle moving with a speed very nearly equal to that of light. Suppose that a particle of rest mass m and total energy \mathcal{E} collides head on with a photon of energy \mathcal{E}_γ . Show that the scattered photon has energy

$$\mathcal{E} \left[1 + \frac{m^2 c^4}{4\mathcal{E}\mathcal{E}_\gamma} \right]^{-1}$$

Ultra-high energy cosmic rays have energies up to 10^{20} eV. How much energy can a cosmic ray proton transfer to a microwave background photon?

Solution:

The collinear collision is pictured below:



The equations of conservation of energy and momentum are

$$\begin{aligned} E + E_\gamma &= E' + E'_\gamma \\ p - p_\gamma &= p' + p'_\gamma \end{aligned}$$

Since the initial and final photon are massless, we have $E_\gamma = cp_\gamma$ and $E'_\gamma = cp'_\gamma$. Let us solve for the final energy and momentum of the proton:

$$\begin{aligned}
E' &= E + E_\gamma - E'_\gamma \\
cp' &= cp - E_\gamma - E'_\gamma
\end{aligned}$$

Thus

$$\begin{aligned}
E'^2 - (cp')^2 &= (E + E_\gamma - E'_\gamma)^2 - (cp - E_\gamma - E'_\gamma)^2 \\
&= E^2 - (cp)^2 + 2E(E_\gamma - E'_\gamma) + 2cp(E_\gamma + E'_\gamma) - 4E_\gamma E'_\gamma
\end{aligned}$$

Let us make use of the energy-momentum relations $E'^2 - (cp')^2 = m^2 c^4 = E^2 - (cp)^2$ to obtain

$$E'_\gamma [E - cp + 2E_\gamma] = E_\gamma [E + cp]$$

In the extreme relativistic limit, we make the replacement $cp \rightarrow E$ in this expression to obtain

$$E'_\gamma = \frac{EE_\gamma}{E_\gamma + m^2 c^4 / (4E)}$$

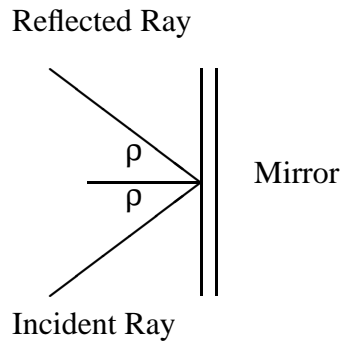
As a crude estimate, that $mc^2 = 10^9 \text{ eV}$ and $E_\gamma = 10^{-4} \text{ eV}$, corresponding to microwave black body radiation at a temperature of a few Kelvins. For an incident cosmic ray proton of energy 10^{20} eV , we obtain $E'_\gamma \approx 4 \times 10^{18} \text{ eV}$. Consequently, the cosmic ray loses 4% of its energy in such a collision. It is felt that this effect is responsible for the decrease in the number of high energy cosmic rays at about this energy.

2. A ray of light is reflected from a plane mirror, which moves in the direction of its normal with velocity v . Prove that the angles of incidence and reflection, θ and ϕ , respectively, are related by the formula

$$\frac{\sin \theta}{\sin \phi} = \frac{\cos \theta \pm v/c}{\cos \phi \mp v/c}$$

Solution:

The incident and reflected rays are shown in the rest frame of the mirror.

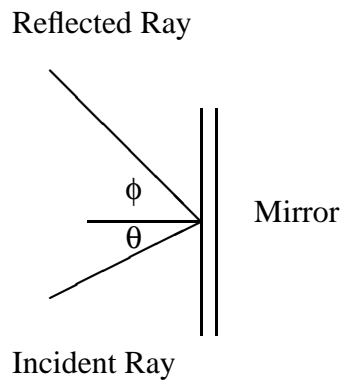


Rest Frame of Mirror

Their four-vector momenta, $(k_i)^\mu$ and $(k_f)^\mu$, are given in that frame as follows:

$$\begin{aligned} (k_i)^\mu &= (k, k \cos \rho, k \sin \rho, 0) \\ (k_f)^\mu &= (k, -k \cos \rho, k \sin \rho, 0) \end{aligned}$$

In the moving frame, these rays are shown below:



Moving Observer

Their four-vector momenta, $(k'_i)^\mu$ and $(k'_f)^\mu$, are given in that frame as follows:

$$\begin{aligned} (k'_i)^\mu &= (k'_i, k'_i \cos \theta, k'_i \sin \theta, 0) \\ (k'_f)^\mu &= (k'_f, -k'_f \cos \rho, k'_f \sin \phi, 0) \end{aligned}$$

The respective four-momenta in the two frames are related by the Lorentz transform. It is convenient to apply the inverse transforms. We do so for the x and y components of the incident four-momentum:

$$\begin{aligned}(k_i)^1 &= \gamma((k'_i)^1 + \beta(k'_i)^0) \\ k \cos \rho &= \gamma k'_i (\cos \theta + \beta)\end{aligned}$$

$$\begin{aligned}(k_i)^2 &= (k'_i)^2 \\ k \sin \rho &= k'_i \sin \theta\end{aligned}$$

Taking the ratio of these equations, we obtain

$$\tan \rho = \frac{\sin \theta}{\gamma (\cos \theta + \beta)}$$

Now we consider the x and y components of the reflected beam:

$$\begin{aligned}(k_f)^1 &= \gamma((k'_f)^1 + \beta(k'_f)^0) \\ -k \cos \rho &= \gamma k'_f (-\cos \phi + \beta)\end{aligned}$$

$$\begin{aligned}(k_f)^2 &= (k'_f)^2 \\ k \sin \rho &= k'_f \sin \phi\end{aligned}$$

Taking the ratio of these equations, we obtain

$$\tan \rho = \frac{\sin \phi}{\gamma (\cos \phi - \beta)}$$

Since the two expressions for $\tan \rho$ must be identical, we obtain

$$\frac{\sin \theta}{\gamma (\cos \theta + \beta)} = \frac{\sin \phi}{\gamma (\cos \phi - \beta)}$$

The factors of γ cancel out, so that

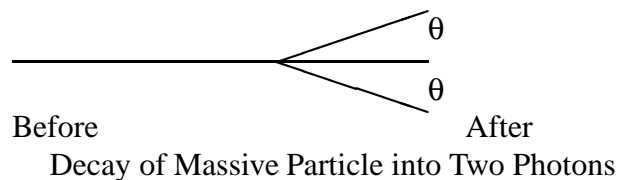
$$\frac{\sin \theta}{\cos \theta + \beta} = \frac{\sin \phi}{\cos \phi - \beta}$$

The result is thus established.

3. A particle of mass m and total energy E decays into two massless particles that travel symmetrically with respect to the direction of motion of the original particle. Determine the angle θ of the decay products with respect to the original direction. As particular cases, consider the non-relativistic case and the extreme relativistic case.

Solution:

Here is a picture of the situation before and after the decay process:



Let the incident particle have initial energy \mathcal{E} and momentum \vec{p} to the right. The final photons have momenta \vec{p}_{\pm} and energies $\mathcal{E}_{\pm} = cp_{\pm}$:

$$\begin{aligned}\mathcal{E} &= \mathcal{E}_+ + \mathcal{E}_- \\ \vec{p} &= \vec{p}_+ + \vec{p}_-\end{aligned}$$

It follows from symmetry that $\mathcal{E}_+ = \mathcal{E}_-$, and that the momenta \vec{p}_{\pm} are of equal magnitude. Thus

$$\begin{aligned}\mathcal{E} &= 2\mathcal{E}_+ \\ p &= 2p_+ \cos \theta\end{aligned}$$

As a consequence, with v the speed of the incident particle, we obtain

$$\frac{v}{c} = \frac{cp}{\mathcal{E}} = \frac{cp_+ \cos \theta}{\mathcal{E}_+} = \cos \theta$$

Note that, when the speed of the incident particle is small compared with the velocity of light, the opening angle between the two photons (2θ) is close to 90° . As the speed of the incident particle increases, that opening angle becomes smaller, approaching zero in the extreme relativistic limit.

4. Calculate the area of a circle of radius r (distance from center to circumference) in the two-dimensional geometry that is a sphere of radius a . Show that this reduces to πr^2 when $r \ll a$.

Solution:

Choose a spherical polar coordinate system with the polar axis (North pole) at the center of the circle. The radius of the sphere is $r = a \theta_0$, where θ_0 is the polar angle subtended by the perimeter of the circle. The surface area inside the (curved) circle is

$$A = a^2 \int_0^{\theta_0} d\theta \sin \theta \int_0^{2\pi} d\phi = 2\pi a^2 (1 - \cos \theta_0) = 2\pi a^2 (1 - \cos \frac{r}{a})$$

From the relation

$$\cos \theta_0 = 1 - \frac{\theta_0^2}{2} + \frac{\theta_0^4}{24} - \dots$$

we obtain

$$A = 2\pi a^2 \left(\frac{\theta_0^2}{2} - \frac{\theta_0^4}{24} \right) = \pi r^2 \left(1 - \frac{\theta_0^2}{12} \right) = \pi r^2 \left(1 - \frac{r^2}{12a^2} \right)$$

When $r = 100 \text{ m}$, the area is approximately $\pi \times 10^4 \text{ m}^2 = \pi$ hectares, differing from that value by about 2×10^{-11} . When $r = 100 \text{ km}$, the deviation is approximately 2×10^{-5} .

5. Consider a particle with four-momentum p and an observer with four-velocity $u = (\gamma c, \gamma \vec{v})$. Show that, if the particle goes through the observer's laboratory, the magnitude of the three-momentum measured is

$$|\vec{p}'|^2 = \left(\frac{p \cdot u}{c} \right)^2 - p \cdot p$$

Solution:

Suppose that the observer is moving in the x -direction with speed v , so that $u^\mu = \gamma(c, v, 0, 0) = \gamma c(1, \beta, 0, 0)$.

The four-momentum of the particle in the initial frame is $p^\mu = (p^0, \vec{p})$, so that

$$\begin{aligned}
p \cdot p = p^\mu p_\mu &= (p^0)^2 - \vec{p} \cdot \vec{p} = p_0^2 - p_x^2 - p_y^2 - p_z^2 \\
p \cdot u = p^\mu u_\mu &= p^0 u^0 - \vec{p} \cdot \vec{u} = p^0 u^0 - p_x u_x = \gamma c (p^0 - \beta p_x)
\end{aligned}$$

We thus obtain

$$\left(\frac{p \cdot u}{c}\right)^2 = \gamma^2 (p^0 - \beta p_x)^2$$

Consequently

$$\begin{aligned}
\left(\frac{p \cdot u}{c}\right)^2 - p \cdot p &= \gamma^2 (p^0 - \beta p_x)^2 - ((p^0)^2 - p_x^2 - p_y^2 - p_z^2) \\
&= (\gamma^2 - 1)(p^0)^2 - 2\gamma^2 \beta p^0 p_x + (1 + \beta^2 \gamma^2) p_x^2 + p_y^2 + p_z^2 \\
&= \gamma^2 (p_x^2 - 2\beta p_x p^0 + \beta^2 (p^0)^2) + p_y^2 + p_z^2 \\
&= \gamma^2 (p_x - \beta p^0)^2 + p_y^2 + p_z^2
\end{aligned}$$

The components of spatial momentum in the frame of the moving observer are

$$\begin{aligned}
p'_x &= \gamma(p_x - \beta p^0) \\
p'_y &= p_y \\
p'_z &= p_z
\end{aligned}$$

so that

$$|\vec{p}'|^2 = \gamma^2 (p_x - \beta p^0)^2 + p_y^2 + p_z^2$$

The result is thus established.